Progress on the theory wishlist for multiloop integrals

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Feynman integrals as iterated integrals (I)

At one loop, only logarithm and dilogarithm needed

$$\log z = \int_{1}^{z} \frac{dt}{t} \qquad \qquad \text{Li}_{2}(z) = \int_{0}^{z} \frac{dt_{1}}{t_{1}} \int_{0}^{t_{1}} \frac{dt_{2}}{1 - t_{2}}$$

- what functions will appear at higher loops?
- Logarithm and dilogarithm are first examples of iterated integrals with special ``d-log`` integration kernels

$$\frac{dt}{t} = d\log t \qquad \frac{-dt}{1-t} = d\log(1-t) \qquad \frac{dt}{1+t} = d\log(1+t)$$

• these are called harmonic polylogarithms (HPL) [Remiddi, Vermaseren]

e.g.
$$H_{1,-1}(x) = \int_0^x \frac{dx_1}{1-x_1} \int_0^{x_1} \frac{dx_2}{1+x_2}$$

- shuffle product algebra
- coproduct structure
- Mathematica implementation [Maitre]
- weight: number of integrations

Feynman integrals as iterated integrals (2)

Natural generalization: multiple polylogarithms

[also called hyperlogarithms; Goncharov polylogarithms]

allow kernels $w = d \log(t - a)$

$$G_{a_1,...a_n}(z) = \int_0^z \frac{dt}{t - a_1} G_{a_2,...,a_n}(t)$$

numerical evaluation: GINAC [Vollinga, Weinzierl]

• Chen iterated integrals

$$\int_C \omega_1 \omega_2 \ldots \omega_n$$
 $C:[0,1] \longrightarrow M$ (space of kinematical variables)

Alphabet: set of differential forms $\omega_i = d \log \alpha_i$

integrals we discuss will be monodromy invariant on $M \setminus S$ S (set of singularities)

more flexible than multiple polylogarithms!

- Uniform weight functions (pure functions):
 - Q-linear combinations of functions of the same weight

d-log representations

Can we make it manifest when integrals evaluate to pure functions?

$$\mathcal{A}_{4}^{\ell=0} \times \bigcap_{p_{1}}^{p_{2}} \mathcal{A}_{4}^{\ell=0} \times \int_{p_{4}}^{\ell=0} \mathcal{A}_{4}^{\ell=0} \times \int_{p_{4}}^{\ell$$

[Caron-Huot, talk at Trento, 2012]

[Lipstein, Mason, 2013]

$$\frac{d^{4}\ell (p_{1} + p_{2})^{2}(p_{1} + p_{3})^{2}}{\ell^{2}(\ell + p_{1})^{2}(\ell + p_{1} + p_{2})^{2}(\ell - p_{4})^{2}}$$

$$= d \log \left(\frac{\ell^{2}}{(\ell - \ell^{*})^{2}}\right) d \log \left(\frac{(\ell + p_{1})^{2}}{(\ell - \ell^{*})^{2}}\right) d \log \left(\frac{(\ell + p_{1} + p_{2})^{2}}{(\ell - \ell^{*})^{2}}\right) d \log \left(\frac{(\ell - p_{4})^{2}}{(\ell - \ell^{*})^{2}}\right)$$

very suggestive! New ways of performing loop integrations?

Cuts and integrated integrands

- use cuts of integrals as guiding principle for finding convenient integral basis [JMH, PRL 110 (2013) 25]
- integrals with simple cuts are expected to integrate to uniform weight functions

idea: any cut that completely localizes the integral should give just a rational number

Strategy for computing Feynman integrals using differential equations

- Useful facts:
 - (I) For a given problem, one can choose a finite basis of Feynman integrals
 - (2) Basis integrals satisfy coupled first-order differential equations
 - (3) many classes of Feynman integrals evaluate to iterated integrals
- Idea: choose basis such that the differential equations are simple,
 and such that (3) is made obvious

Key points of the method

- ullet differential equations for master integrals f
- crucial: choose convenient basis (systematic procedure)
 - \rightarrow makes solution trivial to obtain
- elegant description: Feynman integrals specified by:
 - (1) set of 'letters' (related to singularities x_k)
 - (2) set of constant matrices A_k

Example: one dimensionless variable
$$x$$
; $D=4-2\epsilon$
$$\partial_x \vec{f}(x;\epsilon) = \epsilon \sum_k \frac{A_k}{x-x_k} \vec{f}(x;\epsilon)$$

- ullet expansion to any order in ϵ is linear algebra answer: multiple polylogarithms of uniform weight ('transcendentality')
- asymptotic behavior $\vec{f}(x;\epsilon) \sim (x-x_k)^{\epsilon A_k} \vec{f}_0(\epsilon)$
- natural extension to multi-variable case

Multi-variable case and the alphabet

• Natural generalization to multi-variable case

$$d\vec{f}(\vec{x};\epsilon) = \epsilon d \left[\sum_{k} A_k \log \alpha_k(\vec{x}) \right] \vec{f}(\vec{x};\epsilon)$$
constant matrices letters (alphabet)

Examples of alphabets:

4-point on-shell

$$\alpha = \{x, 1 + x\}$$

two-variable example (from I-loop Bhabha scattering):

$$\alpha = \{x, 1 \pm x, y, 1 \pm y, x + y, 1 + xy\}$$

more complicated examples later

- Matrices and letters determine solution
- Immediate to solve in terms of Chen iterated integrals

Important points differential equations

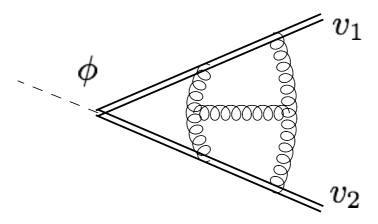
- Uniform weight basis can be found systematically using cuts
 (related to d-log representations)
 [Arkani-Hamed et al.]
 Other ideas
 [Mastrolia et al.]
 [Caron-Huot, J.M.H.]
 [Gehrmann et al.]
- DE provide information about integrals in compact form (alphabet, matrices)
- contain more information than epsilon expansion: exact limits
- boundary conditions often for free (e.g. finiteness in certain limits) application: bootstrap for single-scale integrals [J.M.H., A.V. Smirnov, V.A. Smirnov]

- Chen iterated integrals give most compact form of answer
- To given weight, answer can be rewritten in terms of minimal function basis

3-loop HQET integrals

• 8 integral families, e.g.

$$\cos \phi = \frac{v_1 \cdot v_2}{\sqrt{v_1^2} \sqrt{v_2^2}}, \quad x = e^{i\phi}$$



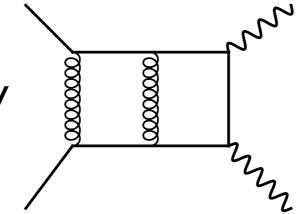
- alphabet
- $\alpha = \{x, 1+x, 1-x\}$
- 71 master integrals
- application: QCD cusp anomalous dimension

[Grozin, J.M.H., Korchemsky, Marquard, to appear 2014]

physics motivation: infrared divergences of massive scattering amplitudes

Vector boson production integrals $pp \rightarrow VV'$

• sample integral family



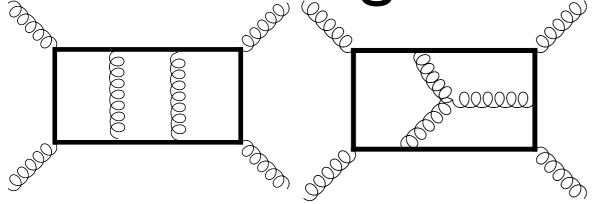
[JMH, Melnikov, V. Smirnov, JHEP 1430 (2014)] [JMH, Caola, Melnikov, V. Smirnov, 1404.5590]

- variables S, T, M_3^2, M_4^2
- physical region 0 < x, 0 < y < z < 1
- alphabet

$$\alpha = \{x, y, z, 1 + x, 1 - y, 1 - z, 1 + xy, z - y, 1 + y(1 + x) - z, xy + z, 1 + x(1 + y - z), 1 + xz, 1 + y - z, z + x(z - y) + xyz, z - y + yz + xyz\}.$$

• boundary condition computed at $x \to 0$, $y \to z \to 1$

Massive integrals for light-by-light scattering



[Caron-Huot, J.M.H., 2014]

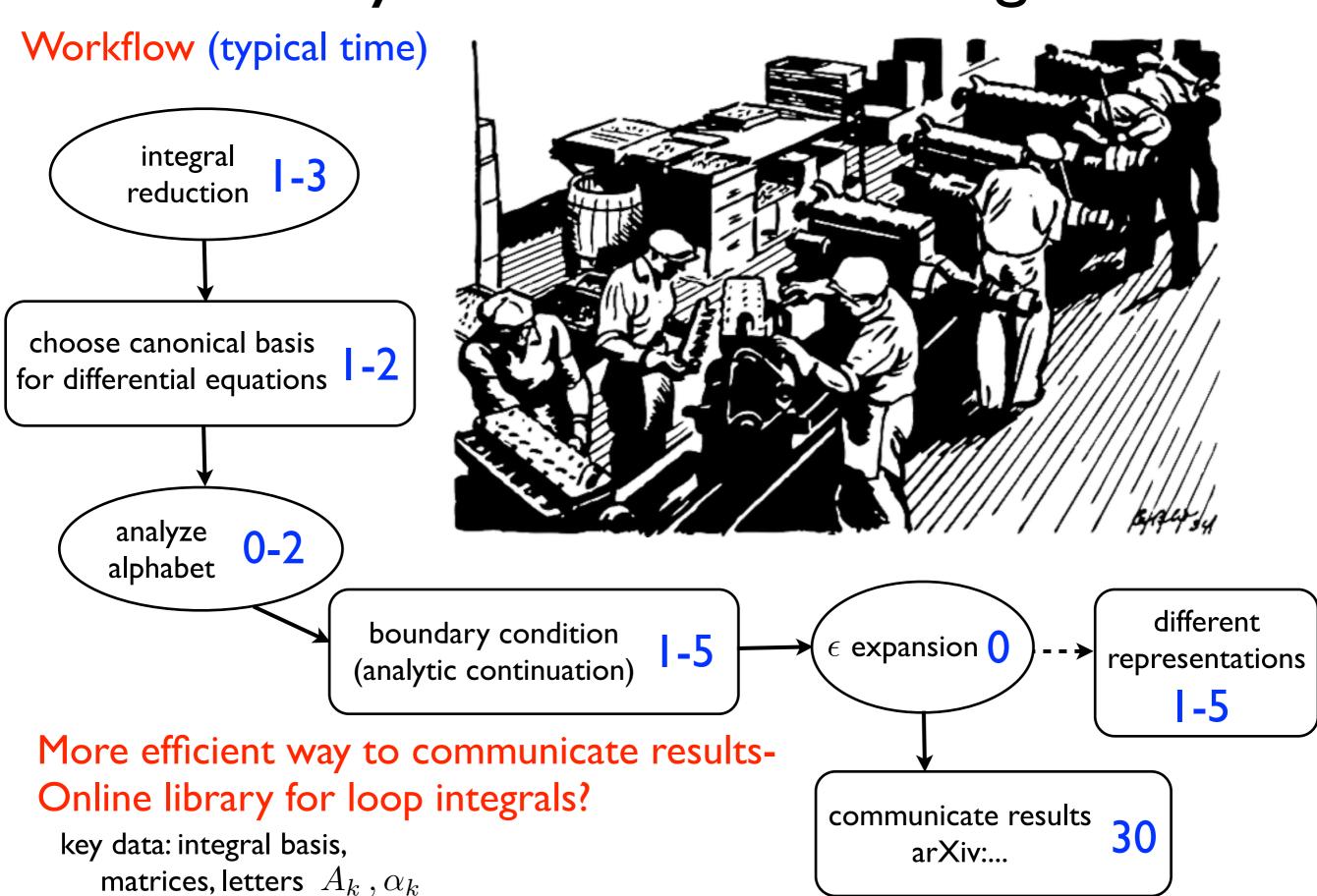
- variables m^2, s, t
 - 3 loops and 3 scales!
- full set of 2-loop master integrals (at 3 loops: all finite master integrals in D=4) similar integrals in QCD, e.g. for finite top quark mass
- alphabet

$$\begin{split} \alpha = & \Big\{ u, 1+u, v, 1+v, u+v, \frac{\beta_u-1}{\beta_u+1}, \frac{\beta_v-1}{\beta_v+1}, \frac{\beta_{uv}-1}{\beta_{uv}+1}, \frac{\beta_{uv}-\beta_u}{\beta_{uv}+\beta_u}, \frac{\beta_{uv}-\beta_v}{\beta_{uv}+\beta_v}, \\ & u^2-4v, v^2-4u, \frac{2-2\beta_{uv}+u}{2+2\beta_{uv}+u}, \frac{2-2\beta_{uv}+v}{2+2\beta_{uv}+v}, \\ & 1+u+v, \frac{4-v+\beta}{4-v-\beta}, \frac{4+v+\beta}{4+v-\beta}, \frac{(4\beta_u+\beta)(4\beta_u+\beta_uv+\beta)}{(4\beta_u-\beta)(4\beta_u+\beta_uv-\beta)}, \frac{(4\beta_{uv}+\beta)(4\beta_{uv}-\beta_{uv}v+\beta)}{(4\beta_{uv}-\beta)(4\beta_{uv}-\beta_{uv}v-\beta)} \Big\} \\ & \text{where } u = -4m^2/s \,, \quad v = -4m^2/t \,, \\ & \beta_u = \sqrt{1+u} \,, \quad \beta_v = \sqrt{1+v} \,, \qquad \beta_{uv} = \sqrt{1+u+v} \,, \qquad \beta = \sqrt{16+16u+8v+v^2} \end{split}$$

• efficient numerical representation for Chen iterated integrals

$$g_{37}(2,4) = 0.0764922717271986970254859257468...$$

Factory-line for master integrals



Thank you!

Extra slides

A word of caution: more exotic objects

- mathematicians like to consider single-scale Feynman integrals
- conjecture that certain periods only evaluate to multiple zeta values (MZV) appear disproven by [Brown, Schnetz]
- Elliptic functions

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relevant e.g. in top quark physics Czakon et al.
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also appear in massless N=4 SYM [Caron-Huot, Larsen]

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recent work Elliptic polylogairthms [Brown, Levin]
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[Bloch, Vanhove] [Vanhove] [Remiddi, Tancredi] [Adams, Bogner, Weinzierl]

Note: weight property generalizes weight n -> (n/2,n/2) mixed Hodge theory systematic and practical way for dealing with them for practical applications?

• Here: cases where Chen iterated integrals are sufficient

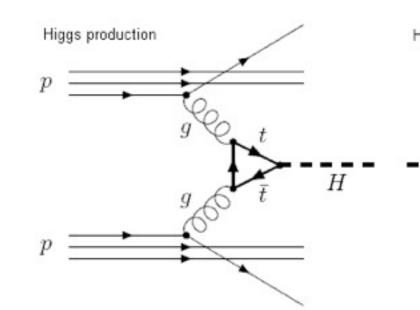
The alphabet and perfect bricks (I)

Can we parametrize variables such that alphabet is rational?

Not essential, but nice feature.

Example: Higgs production

encounter
$$\sqrt{1-4m^2/s}$$
 choose $-m^2/s=x/(1-x)^2$ $\alpha=\{x,1-x,1+x\}$ (to two loops)



Note: this is a purely kinematical question. Independent of basis choice.

Related to diophantine equations
 e.g. find rational solutions to equations such as

$$1 + 4a = b^2$$

here we found a 1-parameter solution

$$a = \frac{x}{(1-x)^2} \qquad b = \frac{1+x}{1-x}$$

The alphabet and perfect bricks (2)

• Classic example: Euler brick problem

Find a brick with sides a,b,c and diagonals d,e,f integers

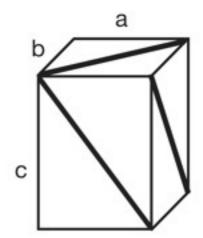
smallest solution (P. Halcke):

$$(a,b,c)=(44,117,240)$$

$$a^2 + b^2 = d^2,$$

$$a^2 + c^2 = e^2$$
,

$$b^2 + c^2 = f^2.$$



Perfect cuboid (add eq. $a^2 + b^2 + c^2 = g^2$): open problem in mathematics!

Similar equations for particle kinematics

e.g encountered in 4-d light-by-light scattering

$$u = -4m^2/s \qquad v = -4m^2/t$$

$$\beta_u = \sqrt{1+u}, \ \beta_v = \sqrt{1+v}, \ \beta_{uv} = \sqrt{1+u+v}$$

Need two-parameter solution to

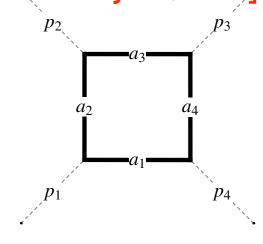
$$\beta_u^2 + \beta_v^2 = \beta_{uv}^2 + 1$$

e.g.
$$\beta_u=rac{1-wz}{w-z}\,, \ \ \beta_v=rac{w+z}{w-z}\,, \ \ \beta_{uv}=rac{1+wz}{w-z}\,.$$

more roots in D-dim and at 3 loops! - in general alphabet changes with the loop order!

Find such solutions systematically? Minimal polynomial order?

[Caron-Huot JMH, 2014]



Goncharov weight four conjecture

rewrite any multiple polylogarithm in terms of function basis

[Goncharov]

e.g. at weight 4 (important for NNLO computations)

$$\{\log(x)\log(y)\log(z)\log(w), \log(x)\log(y)\operatorname{Li}_{2}(z), \operatorname{Li}_{2}(x)\operatorname{Li}_{2}(y), \log(x)\operatorname{Li}_{3}(y), \operatorname{Li}_{4}(x), \operatorname{Li}_{2,2}(x,y)\}$$

for set of arguments (to be found - symbol/coproduct provides guidance)

minimal set of integration kernels vs. minimal set of function arguments

practical tool: `symbol`` useful projections [Goncharov, Spradlin, Vergu, Volovich]
 [Brown] [Goncharov]

e.g. project on $\mathrm{Li}_{2,2}(x,y)$ part e.g. project out all products

lecture notes: [Vergu] [Brown][Zhao]

• ``symbol`` = Chen iterated integral without boundary information diff. eqs. or other information can be used to fix this

Equivalent representations

• version I: Chen iterated integrals

$$g_6 = \int_{\gamma} d \log \frac{\beta_u - 1}{\beta_u + 1} d \log \frac{\beta_{uv} - \beta_u}{\beta_{uv} + \beta_u} + \int_{\gamma} d \log \frac{\beta_v - 1}{\beta_v + 1} d \log \frac{\beta_{uv} - \beta_v}{\beta_{uv} + \beta_v}.$$

[2 loops: 10 terms]

version 2: Goncharov polylogarithms

(if alphabet rational in at least one variable)

$$g_6 = -G_{-1,0}(w) + G_{0,-1}(w) - G_{0,1}(w) + G_{1,0}(w) + H_{-1,0}(z) - H_{0,-1}(z) - H_{0,1}(z) + H_{1,0}(z) - G_0(w)H_{-1}(z) + G_{-1}(w)H_0(z) - G_1(w)H_0(z) - G_0(w)H_1(z).$$

[2 loops: 2-3 pages]

• version 3: minimal function basis $g_6 = -\beta_{uv}/2I_1$

$$I_{1} = \frac{2}{\beta_{uv}} \left\{ 2 \log^{2} \left(\frac{\beta_{uv} + \beta_{u}}{\beta_{uv} + \beta_{v}} \right) + \log \left(\frac{\beta_{uv} - \beta_{u}}{\beta_{uv} + \beta_{u}} \right) \log \left(\frac{\beta_{uv} - \beta_{v}}{\beta_{uv} + \beta_{v}} \right) - \frac{\pi^{2}}{2} + \sum_{i=1,2} \left[2 \operatorname{Li}_{2} \left(\frac{\beta_{i} - 1}{\beta_{uv} + \beta_{i}} \right) - 2 \operatorname{Li}_{2} \left(-\frac{\beta_{uv} - \beta_{i}}{\beta_{i} + 1} \right) - \log^{2} \left(\frac{\beta_{i} + 1}{\beta_{uv} + \beta_{i}} \right) \right] \right\}.$$

[2 loops: several pages]

 $\beta_u = \sqrt{1+u}, \ \beta_v = \sqrt{1+v}, \ \beta_{uv} = \sqrt{1+u+v}$

• some examples from literature: [Goncharov et al.] [Duhr] [Gehrmann et al.] ...

[most compact]

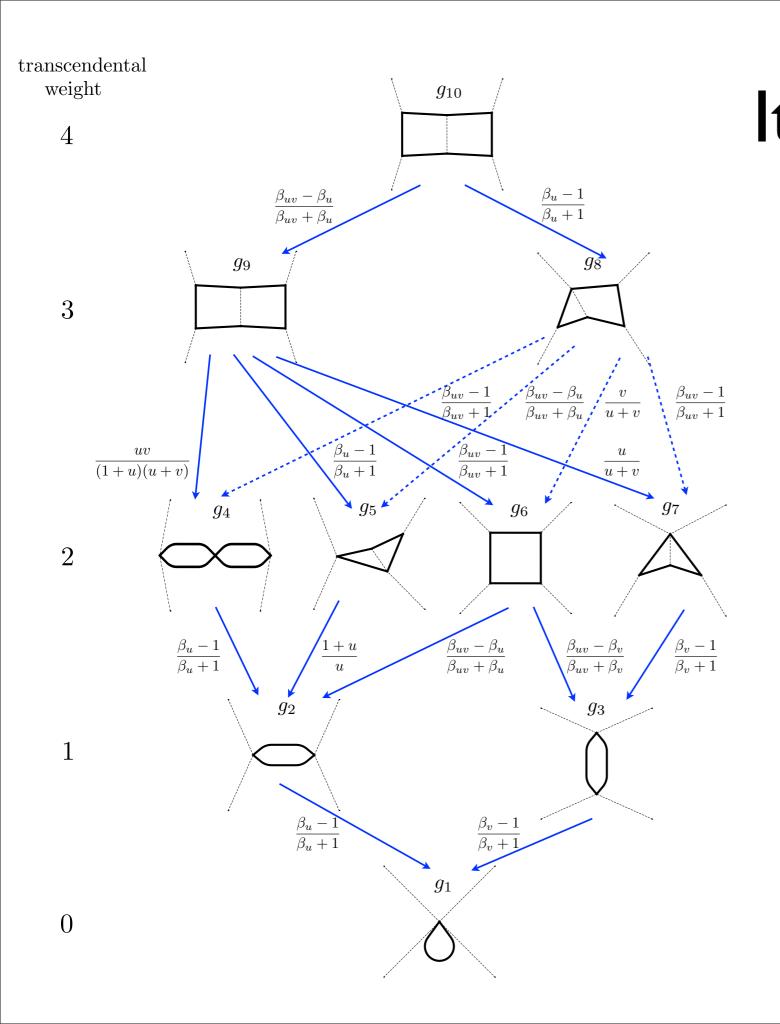
[flexible: analytic continuation, limits]

[easy to see DE, cuts]

[ideas for numerics: J.M.H., Caron-Huot]

[longer expressions; requires rational alphabet; GINAC numerical evaluation]

[arbitraryness; usually long expressions; good at low weight; fast numerical evaluation]



Iterative structure for finite loop integrals

[Caron-Huot, J.M.H. (2014)

- block triangular matrix structure (weight grading)
- algorithm for finding this form